High-Resolution He I and He II Photoelectron Spectra of TiCl₄, SnCl₄, and $(CH_3)_4\$ Sn

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Received March 4, 1982

High-resolution He I and He II photoelectron spectra of TiCl₄, SnCl₄, and (CH₃)₄Sn have been reported. A number of new peaks have been resolved in these spectra which have led to a reassessment of previous TiCl₄ assignments. We suggest that the ordering of the three high binding valence levels in TiCl₄ should be (1e ~ 2a₁) < 3t $2a_1$, as proposed earlier. This order is consistent with theoretical SCF-X α predictions, cross-section and vibrational calculations, and intensity changes in the He I and He II spectra of $SnCl₄$ and $(CH₃)₄Sn$.

Introduction

Ultraviolet photoelectron spectroscopy (UPS) has proven to be of central importance in elucidating the molecular orbital sequences in a wide range of inorganic molecules (for example, see the recent review on transition metal compounds³). In many cases, however, the valence-band spectra are very complex, and the assignments are difficult to make. In the last few years, molecular orbital calculations combined with theoretical cross-section changes between He I α (21.2 eV) and He II α (40.8 eV) spectra⁴⁻⁸ have aided the assignments.

Despite the availability of high-quality molecular orbital calculations and He **I** and He I1 spectra, the assignment of many spectra in the literature is controversial. A good example of the assignment problems is provided by the UPS spectra of the relatively simple but fundamentally important molecule TiC14. Experimental He I spectra have been published five times, $5.9-12$ the He II spectrum has been reported,⁵ and a large number of molecular orbital calculations have been published, including INDO,¹³ SCF-X α ,^{11,14,15} and ab initio¹⁶ calculations. Despite all of this effort, Cowley states³ that some of "the bands in the $TiCl₄$ spectrum cannot be assigned unequivocably" because the calculations differ in the sequencing of molecular orbitals and the UPS bands are quite closely spaced.

All of the above $TiCl₄$ spectra, and most of the spectra of inorganic molecules in the literature, have been recorded at medium resolutions, usually about *60* meV for He I spectra and 100 meV for He I1 spectra. In our studies of ligand field splittings and valence-band spectra, we have observed resolutions for *both* He I and He I1 spectra in the 30-50-meV region.^{7,17,18} In this paper, we use this high-resolution com-

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bined with theoretical cross-section and vibronic coupling calculations to show up several additional features in the $TiCl₄$ spectra which lead to a reassessment of the spectral assignments. We also report selected spectra of $SnCl₄$ and $(CH₃)₄Sn$ to enable us to understand better the intensity changes between He I and He II $Ticl_4$ valence-band spectra.

Experimental Section

(A) Ultraviolet **Photoelectron Spectra** High-punty TiC1, and SnCI, were purchased from Fisher Scientific. SnCl₄ was distilled prior to recording its photoelectron spectra. $(CH₃)₄Sn$ was prepared according to the method described by Edgell and Ward.¹⁹ The purity of this compound was confirmed by NMR.

All photoelectron spectra were measured in a McPherson ESCA 36 spectrometer at room temperature with use of a hollow cathode lamp described previously.¹⁷ The spectra were obtained by multiscanning the region of interest. We routinely obtain Ar 3p_{3/2} line widths of 17 and 26 meV with He **I** and He **I1** radiation, respectively. TiC1, proved to be a very difficult compound to work with. After spectra were run for a short time, it destroyed our channeltron detectors. The statistics on our He **I1** spectrum are not as good as those on our normal He **I1** spectra because of this problem.

The photoelectron **peaks** were fitted with a Voigt function simulated by a linear combination of Gaussian-Lorentzian line shapes.²⁰ Four parameters are used to describe each photoelectron peak: the peak position (binding energy), the full width at half-maximum (fwhm), the peak height (YMAX), and the Gauss fraction (GFAC), which is a measure of the Gaussian character of the peak. The "goodness of fit" criterion requires that the sum of the squares of the deviation from the best fit divided by the variance of a single count is minimized. From the fitting procedure, we obtain the best estimates (and standard deviations) for the four peak parameters.

(B) Theoretical Calculations. To confirm the proposed assignments for $TiCl₄$ and $SnCl₄$, we have performed theoretical photoionization cross-section calculations for all the valence molecular orbitals as a function of the photon energy using the X_{α} -SW method.^{21,22} In this method, the continuum wave functions for the photoelectrons are generated from the ground-state $X\alpha$ -SW molecular potential. Therefore, the continuum orbitals possess the correct molecular symmetry and orthogonality to the ground-state wave functions.²³ All symmetry-allowed photoionization processes (dipolar selection rule) are included in the calculation. From previous experience, 24.25 we learned that for big molecules, the absolute magnitude of the theoretical photoionization cross section is **very** sensitive to the completeness of the partial wave expansion of the continuum wave functions. The merit of the X_{α} -SW method in predicting photoionization cross sections for molecules has been established in several recent studies.²⁴⁻²⁷

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Figure 1. He **I** and He **I1** photoelectron spectra of TiC14.

In order to establish the reliability of the present theoretical results, we have performed calculations with the maximum azimuthal quantum number for the continuum wave functions expanded to *I* = **5** and *¹* = **7.** Tangential spheres were used in all calculations.

We have obtained qualitative agreement for the variation of the photoionization cross sections with photon energy from both calculations. We observed that there are moderate variations in the absolute magnitude of the cross sections. However, this should not affect the qualitative picture below.

We adopted the **Green's** function method to calculate the vibrational structure for the molecular ions.^{28,29} The renormalized ionization energies and vibronic coupling parameters are calculated by the diagonal 2ph-TDA (two-particle-hole Tamm-Dancoff approximation) procedure.29 The relevant ground-state wave functions and various Coulomb and exchange integrals are obtained from an INDO/2S program parameterized for transition metal compounds.^{30,31}

Results and Discussion

(A) Valence-Band Spectra and Assignment. The He I and He II photoelectron spectra of the entire $TiCl₄$ valence-band region are shown in Figure 1, while narrow energy scans of small regions of the He I spectrum are shown in Figure 2. All peak parameters for these spectra are given in Table I. The gross features of our spectra compare well with the spectra reported previously. $5,9-12$

Using Orchard's⁵ band notation, we find four well-resolved bands at 11.7 eV **(A),** 12.8 eV (B), 13.2 eV (C + D), and 13.9 eV (E). Closer examination of our spectra in Figures 1 and 2 reveals that every one of these bands is composed of two or more peaks. The statistics in the He I1 spectrum are not good enough to resolve the finer details in the bands **A,** B, and E. Our instrumental resolution is indicated by the 30-meV line widths of the component peaks of band E (Figure 2C and Table I). None of the previously published spectra of $TiCl₄$ showed even any signs of asymmetry on this band.

In the tetrahedral $TiCl₄$ molecule, the metal-halogen bonding molecular orbitals transform as $e + t_2$ symmetries, and the halogen lone pairs, as $t_1 + t_2 + e + a_1$ symmetries. The assignment of bands A and B to the $1t_1$ and $3t_2$ orbitals, respectively, of largely Cl 3p character is generally agreed on,^{3,5} and our spectra are consistent with this assignment. Both bands **A** and B show very distinct shoulders, and these bands have both been fitted to two peaks-a and a' separated by 146 meV and b and b' separated by 190 meV (Table I). These distinct shoulders, giving these large splittings, preclude the

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Figure 2. Expanded He I photoelectron spectra of TiC14: **(A)** band **A;** (B) bands B, C, and D; (C) band E.

two doublets being due solely to vibrational structure. Other possible causes are spin-orbit splitting and Jahn-Teller splitting.^{11,12} Both Cl 3p spin-orbit splittings (\leq 73 meV) and Jahn-Teller splittings are normally considerably smaller than our observed splitting, $11,12,32$ but both are well-known in tetrahedral AB₄ molecules.^{11,12} The overall observed structure is probably due to a combination of vibrational, spin-orbit, and Jahn-Teller splittings. The vibrational broadening is calculated below.

The assignment of the next three peaks has been in dispute. Orchard et al. proposed two peaks in the $C + D$ band region and assigned them to the 1e and $2t_2$ orbitals without specifying the order. They assigned band E to the $2a_1$ orbital.¹⁰⁻¹² Combined with molecular orbital calculations, our spectra strongly suggest (but do not prove unequivocally) that the order of levels is not le, $2t_2 < 2a_1$ but rather (le $\sim 2a_1$) < $2t_2$.

Figures 1 and 2B show that the $C + D$ region contains two peaks separated by 0.29 eV. The four-peak structure in band E resembles a normal vibrational progression. We associate

Table 11. Calculated MO Charge Distribution of TiC1, from the *Xa-SW* **Method**

		% charge distribution					
			Ti				
band	orbital	đ	S	р	CI	outsphere	intersphere
А	lt,		.		81		
В	3t,				76		21
C	1e	12			66		21
Đ	2a ₁				75		14
E	2t ₂	25			63		

the 33-meV energy difference (266 cm⁻¹) with the ν_1 fundamental TiCl₄ vibration of 388 cm⁻¹.³³ This substantial reduction observed in the ν_1 energy (31%) is extremely similar to that observed on the corresponding highest binding energy band in the $SiCl₄ spectrum.⁶$ In both cases, this structure must arise from a bonding M-Cl molecular orbital.

The relative intensities of the He I and He I1 peaks are crucial in the interpretation. The intensity of bands **A,** B, and D diminish relative to C on going from He I to He 11, while band E increases sharply. It has often been shown that the C1 3p cross section decreases from He I to He 11, while the metal d orbital cross section increases.³⁴ Following this empirical "rule", we would expect bands **A,** B, and D to be associated with mostly C1 3p character, with bands C and E being associated with orbitals with substantial Ti 3d character.

Before assigning bands C, D, and **E** with the above intensity and vibrational information, it is important to examine the latest ab initio and $X\alpha$ calculations. The ab initio calculation¹⁶ assigns these last three orbitals to $1e < 2t_2 < 2a_1$, with 16% , 21%, and 9% Ti 3d character, respectively. This assignment agrees with that of Orchard et al.⁵ The X α calculations^{11,15} give the order $1e < 2a_1 < 2t_2$ but do not give orbital characters. To obtain these orbital characters, we have performed an SCF-X α -SW calculation³⁵ on TiCl₄. The relative ordering of our ionization energies calculated by the transition-state procedure³⁶ agrees with the results of the previous X_{α} calculations and our INDO calculation described below. The charge distribution from the calculation is given in Table 11. The results agree qualitatively with those from the ab initio calculations, with the Ti 3d character in these three orbitals decreasing in the order $2t_2 > 1e > 2a_1^{37,38}$

The X_{α} assignment (1e < 2a₁ < 2t₂) is also consistent with the variation of He I and He I1 intensities, the vibrational structure on band E, and the order of increasing Ti 3d character in these orbitals from both X_{α} and ab initio calculations. The two molecular orbitals with highest Ti 3d character $(2t₂,$ le) increase in intensity the most from He I to He 11, while the C1 3p orbitals decrease the most. The orbital with the largest bonding character $(2t₂)$ gives rise to the vibrational structure with a vibrational spacing ν_1 much reduced from ν_1 for the neutral molecule.

It is interesting to note that the intensity ratio between band C and D does not change significantly from He I (1.13) to He I1 (1.32) radiation. **A** possible explanation for this observation is that vibrational and Jahn-Teller structure on the le orbital (band C) may accidentally overlap with the 2a,

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Table **111.** SnC1, Valence-Band **(1 1-13** eV) He **I** and He **I1** Parameters (eV)

He I α			He IIa					
peak	ΔE^a	fwhm ^a	GFAC ^a	area ^o	ΔE^a	fwhm ^a	GFAC ^a	area ^o
A	0.0	0.34(1)	0.93(6)	43.4	0.0	0.32(3)	0.69(22)	30.4
b	0.23(1)	0.14(1)	1.0	15.1	0.25(2)	0.14(2)	1.0	20.2
b	0.41(1)	0.12(1)	0.62(21)	13.5	0.40(2)	0.12(2)	0.60	15.9
С	0.64(1)	0.18(1)	0.73(2)	28.1	0.64(1)	0.19(1)	0.70	21.8

a Errors given in parentheses. **b** Areas given in arbitrary units.

orbital forming band D. We cannot pinpoint, then, the $2a_1$ peak in these spectra.

Other more subtle intensity changes are apparent in the TiC1, spectrum, and these are not as readily rationalized. Bands A, B, and D are almost entirely C13p in character from both ab initio and $X\alpha$ calculations, and yet the area ratio of A:B:D varies substantially between He I and He I1 spectra (Table I). Could this substantial change in ratio be due to the very small amount of Ti 4p character **[2%** (Table 11) or $4\%^{16}$] in the 3t₂ orbital or Ti 4s character in the 2a₁ orbital, which gives rise to bands B and D, respectively? To answer this question and to further confirm our assignment of the $2a_1$ and $3t_2$ orbitals in TiCl₄, we recorded the He I and He II spectra of the low binding energy region of SnCl₄ and $(CH₃)₄Sn.$

The SnCl₄ spectra (11-13-eV region) are shown in Figure 3 and the peak parameters are summarized in Table 111. Although previous spectra did not resolve peaks b and b' , α our assignment of these bands is the same as that given previously for SnCl_4 .^{6,10} We assign band A to the chlorine 3p nonbonding It, orbital, band B to the Jahn-Teller (or spin-orbit) split slightly bonding $3t_2$ orbital, and band C to the chlorine 3p nonbonding le orbital. This is the same orbital ordering as in TiCl₄, but here there is no possibility of d-orbital involvement in the $SnCl₄$ bonding. The nonbonding natures of the $1t₁$ and 1e orbitals are ensured by symmetry, while our $X\alpha$ -SW calculation on $SnCl₄$ shows that as for TiCl₄, the 3t₂ orbital **possess** a very **small** amount (1.6%) of metal (Sn 5p) character.

It is particularly interesting that the area ratio of C:A remains constant at ~ 0.7 for SnCl₄ for both He I and He II spectra, whereas for TiCl₄ this ratio changed from \sim 0.6 to \sim 2.5 from He I to He II (Table I). These ratios are consistent with the qualitative cross-section arguments given previously. The $1t_1$ and 1e orbitals in SnCl₄ are entirely Cl 3p in character, and we would expect the ratio C:A to remain constant. In contrast, the 1e orbital in $TiCl₄$ has some Ti 3d character (Table 111) and band C increased relative to band A on going from the He I to He I1 spectra. This difference in the relative cross section of band C between $TiCl₄$ and $SnCl₄$ further substantiates our assignment of band E to the $2t_2$ orbital in TiCl,, which has large 3d character. The large d character in the $2t_2$ orbital should result in a large increase in band E from He I to He I1 as is observed.

Figure 3 and Table III show that, as for $TiCl₄$, band B in SnC1, increases in intensity relative to band A from He I and He I1 spectra. The ratio A:B varies from 1.5 to 0.8 for the He I and He I1 spectra, respectively. Again a very small amount of metal p character (1.6% Sn 5p) has apparently resulted in a very noticeable relative increase in intensity from He I to He 11.

To obtain a better indication of the photon energy dependence of metal valence p orbitals, we recorded both He I and He II spectra of $(CH₃)₄$ Sn. Our He I spectrum is very similar to the one recorded previously³⁸ and consists of two very broad peaks at 9.7 eV (X) and 13.4 eV (Y) . To the best of our knowledge, the He I1 spectrum has not been published previously. The peak at **9.7** eV has been assigned to the Sn-C $3t_2$ orbital,³⁸ while the 13.4-eV peak is assigned to C-H molecular orbitals. The area ratio *X:Y* increases sharply from

Figure 3. He **I** and He **I1** photoelectron spectra of SnC1, in the 11-13-eV binding energy region.

Table **IV.** Photoionization Cross-Section Calculations for TiCl, and SnCl, at He I and He **I1** Photon Energy (in Mb)

	TiCl _a		SnCl _a		
	He I	He II	He I	He II	
1t,	22.0	0.04	31.0	0.05	
3t ₂	41.0	0.28	29.5	1.60	
1e	17.1	0.80	19.9	0.55	
2a,	5.36	0.015	3.5	0.17	
2t,	2.11	6.69	18.2	1.02	

He I (0.22) to He II (0.46) . It is well-known that M-C orbitals have $30-40\%$ metal p character,^{39,40} and so we expect that the Sn 5p character in the $3t_2$ orbital will also lie in this range. Since the C-H cross sections decrease by less than a factor of 2 from He I to He II,⁷ the Sn 5p cross section must increase on going from He I to He I1 spectra.

We turn again to the A:B ratio in the tetrachlorides. It is well-known that the C13p nonbonding cross section decreases much more rapidly than the C-H cross section on going from He I to He II.^{5,41} A good example of this behavior is given in the spectra of (CH_3) ₃SnCl⁴¹ It is also important to note the dramatic decrease in the neighboring Ar 3p cross section from He I to He II by a factor of about $10⁴²$ With this drastic decrease in Cl 3p cross section, and an increasing metal p cross section, a relatively small metal p character $(\sim 5\%)$ in the 3t₂ orbital could lead to the changes observed in the $SnCl₄$ and TiC1, **A:B** ratio.

(B) Photoionization Cross-Section and Vibrational Calculations. To further our understanding of photoionization

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Figure 4. Calculated cross sections (Mb) for the valence molecular orbitals of TiC14 as a function of photon energy. The high photon energy region (34-52 eV) is shown in an expanded scale in the insert.

Figure 5. Calculated cross sections (Mb) **for** the valence molecular orbitals of SnC1, as a function of photon energy. The high photon energy region (30-52 eV) is shown in an expanded scale in the insert.

cross-section variations, we undertook the theoretical crosssection calculations on $TiCl₄$ and $SnCl₄$. The results of these calculations are summarized in Figures 4 and *5* (for *1* = 7 with an energy increment of 0.17 Ry). The specific cross sections at the He I and He I1 photon energies are given in Table IV.

From Tables I, 111, and IV one can see that qualitatively there is good agreement between the theoretical and experimental results. In the case of $TiCl₄$, the cross section of the Cl p type orbitals $(1t_1, 3t_2,$ and $2a_1)$ indeed diminish relative to le on going from He I to He II, while orbital $2t_2$, which has 25% d character, increases sharply. This supports our assignment for band E as the 2t₂ orbital and not as $2a_1$.¹⁰⁻¹² The calculations also reflect the variation in the area ratio of A:B:D between He I and He 11. It seems that a small amount of Ti 4p character in the $3t_2$ orbital enhances the cross section relative to the $1t_1$ orbital, which cannot have any metal character. If the area of band D is dominated by the vibrational and Jahn-Teller splitting of the le orbital as we have suggested, it is also clear from the cross-section data why band D is more intense than bands A and **B** in the He I1 spectrum.

In the case of SnCl₄, even though the cross-section ratio le:1t₁ (C:A) changes from He I (~0.65) to He II (~11) whereas experimentally it remains almost a constant (~ 0.7) , the calculations do show that a small amount of Sn 5p character in the $3t_2$ orbital enhance the cross section relative to the $1t_1$ and 1e orbitals, which are pure Cl p type orbitals.

Figures **4** and *5* show shape resonance effects on some of the molecular orbitals at about 10 eV above threshold, with the strongest effect on the $3t_2$ orbital. These shape resonances

Table V. Calculated Koopmans and Renormalized Ionization Energies (eV) and Renormalized Electron-Vibration Coupling Constants for TiC1,

orbital	$\epsilon^{\mathbf{R}}$ (2ph-TDA) $-\epsilon_i$		$_{\sigma}$ R		
Ιt,	15.21	12.53	1.18		
3t,	16.14	13.41	1.72		
1a,	16.17	14.02	1.37		
1e	16.57	14.02	3.21		
2t,	16.94	14.84	2.54		

are the result of the interaction between the photoelectron and the molecular field in the final state. This kind of "final-state effect" has been reported in diatomic molecules such as N_2^{22} and NO,²⁷ and in polyatomic molecules such as SF₆.²⁶ Elaboration of these cross-section results and others on other tetrahalides will be the subject of an ensuing publication. 43 Obviously, experimental verification of these dramatic cross-section changes using synchrotron radiation should provide unambiguous assignments in these and other molecules.

To investigate the electron-vibration coupling effect in the molecular ionization spectra and to further confirm our assignments, we have computed the vibrational spectra for the valence orbitals of TiCl₄. Employing the new INDO/2S parameterization for transition metals, 30,31 we obtained the same ground-state valence molecular orbital sequence as predicted by $X\alpha$ -SW calculations.¹⁵ Including many body effects (electronic relaxation and correlation effects) through the diagonal 2ph-TDA procedure did not alter the orbital ordering but improved the numerical agreement of the ionization potentials with experiment considerably (Table V).

Within the adiabatic approximation,⁴⁴ the Green's function method 28,29 can be extended to include vibrational effects using only data for the ground state. It can be shown²⁹ that nondegenerate electronic states can only couple to totally symmetric vibrational modes in the first-order approximation. The adiabatic approximation may not be valid for degenerate electron states unless the energy difference between electron states is large compared to the vibration spacing (weak Jahn-Teller effect), and this may well not be true in our case discussed above.

There is only one totally symmetric vibrational mode for tetrahedral molecules. If we further assume no frequency change occurs due to the ionization process, the vibrational intensities follows a Poisson distribution

$$
I_i(\epsilon) = \exp(-a_{ij}^R) \sum_n \frac{(a^R)^n}{n!} \delta(\epsilon - \epsilon^R)^i + n\omega)
$$
 (1)

where a_{i}^{R} is the coupling parameter for the *i*th orbital. ϵ_{i} is the renormalized ionization energy and $n\omega$ is the vibrational spacing. The coupling parameter can be calculated from

$$
a_{i}^{R} = \left(\frac{k_{i}}{\omega}\right)^{2} \qquad k_{i} = -\frac{1}{2^{1/2}} \left. \frac{\partial \epsilon^{R_{i}}}{\partial Q} \right|_{Q_{0}}
$$

where Q is the normal coordinate belonging to the A_1 symmetry representation.

Vibration-coupling constants from IND0/2S calculations for TiC14 are tabulated in Table V. The ground-state totally symmetric vibration of 388 cm⁻¹ was used throughout.³³ To mimic the limited instrumental resolution of the photoelectron spectra, we have convoluted each theoretical peak with a Lorentzian function of width 0.05 eV. The theoretical vibrational spectrum of TiCl, is depicted in Figure *6.* The

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Figure 6. Theoretical vibrational spectrum of TiCl₄. The intensity of each peak has been weighted according to the orbital degeneracy.

magnitude of the coupling constant indicates strong vibronic coupling accompanying photoionization for all valence molecular orbitals. The relative line widths of the first three bands are qualitatively produced by the calculations. Thus the calculated full widths at half-height for the three bands A, B, $(C + D)$ are 0.18, 0.21, and 0.30 eV. The vibrational coupling constant for the le orbital is the largest, and the resulting broad line partially explains the large line width for $C + D$. The shoulders are of course not reproduced because of the neglect of Jahn-Teller and spin-orbit splitting. Taken together with all the other experimental and theoretical evidence in this paper, our assignments must be considered relatively reliable.

Conclusions

For the relatively simple inorganic molecule $TiCl₄$, we have obtained high-resolution He I and He **I1** spectra. In combination with X_{α} , INDO, cross-section, and vibrational calculations, we have shown that the previous assignment is not correct. All our evidence points to the assignment $1t_1 < 3t_2$ $(1e, 2a_1)$ $(2t_2, 1e)$ However, even with this large amount of experimental and theoretical evidence, we obviously still do not completely understand aspects of the spectrum-the shoulders on bands A, B, and C and the exact position of the $2a_1$ orbital. Valence-band spectra taken with synchrotron radiation along with more elaborate calculations are required to complete the assignments, even on this simple molecule.

Acknowledgment. We wish to thank the generosity of Professor Zerner and Dr. Davenport in furnishing the computer programs used in the present work. The financial support of the NSERC (Canada) is gratefully acknowledged. E.P. wishes to thank Bar Ilan University (Ramat-Gan, Israel) for financial support.

Studies of the Pfeiffer Effect Developed in Tris(pyridine-2,6-dicarboxylato) terbate(II1) by Phenylalkylamines, Phenylalkylamino Alcohols, and Phenylalkylamino Acids

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Received November 10, 1981

Optical activity induced in the title complex through outer-sphere complexation with a series of structurally related environment substances has been studied by means of circularly polarized luminescence (CPL) spectroscopy. The substrates were found to fall into four classes, which were determined by trends in the induced optical activity and by the nature of the functional groups on the substrates. For certain of the chiral materials, the observed optical activity was pH independent, but for others, the CPL spectra were found to invert upon ionization of the protonated ammonium group. Systematic variation of functionality on the chiral substrates enables one to conclude that ionic, hydrophobic, and hydrogen-bonding mechanisms can all contribute to the association of complex and environment substance.

Introduction

The development of optical activity in a racemic mixture of a labile metal complex upon addition of certain chiral compounds (termed the environment substance) is known as the Pfeiffer effect^{2,3} and is usually thought to arise from a displacement of the equilibrium existing between the enantiomers of the racemic metal complex. While several theories have been advanced to account for the effect, 4 most of the reported work has been explained by the existence of complexation between the racemic complex and the environment substance as leading to the observed optical activity. Several attempts have been made to use the Pfeiffer effect as a means to predict absolute configurations of a variety of substances, $5-7$

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but these correlations have been very limited in their scope.

We have recently been using Pfeiffer optical activity to study f-f optical activity in trigonal lanthanide complexes, since these complexes are too labile to be resolved by ordinary physical or chemical means. Our work has centered on the tris lanthanide complexes of **pyridine-2,6-dicarboxylic** acid (DPA, or dipicolinic acid), as this ligand is known to form very strong complexes with all members of the lanthanide series⁸ and the tris complex is known to possess an approximate *D,* symmetry in fluid solution at room temperature. 9 The first studies were involved with determinations of conditions under which the effect could be observed, and it was found that the $Tb(DPA)$,³⁻ complex could associate with chiral environment substances through both ionic (resolved **tris(ethy1enediamine)chromium-** $(III)^{10}$) and hydrophobic (L-ascorbic acid¹¹) mechanisms.

The studies involving hydrogen bonding of environment substances to aromatic portions of the DPA ligands were found to be of extreme interest, and we have recently studied the pH

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